

Theoretical Modeling of Nanoscale Confined Light: From Metal Nanoparticles to Nanoholes

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Collaborators

Argonne Experimental Groups:

G. Wiederrecht and coworkers,

L. Yin, U. Welp, V. Vlasko-Vlasov and coworkers

Theorists:

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Outline

- Motivation
- Methods
- Results:
 - Isolated Ag Nanowires
 - Nanowire Arrays
 - TIR Excitation
 - Isolated Holes
- Concluding Remarks





Motivation

The nature of light interaction with metal nanoparticle (MNP) or nanohole systems could lead to:

- Nanoscale optical/opto-electronic devices
- Novel chemical and biological sensors

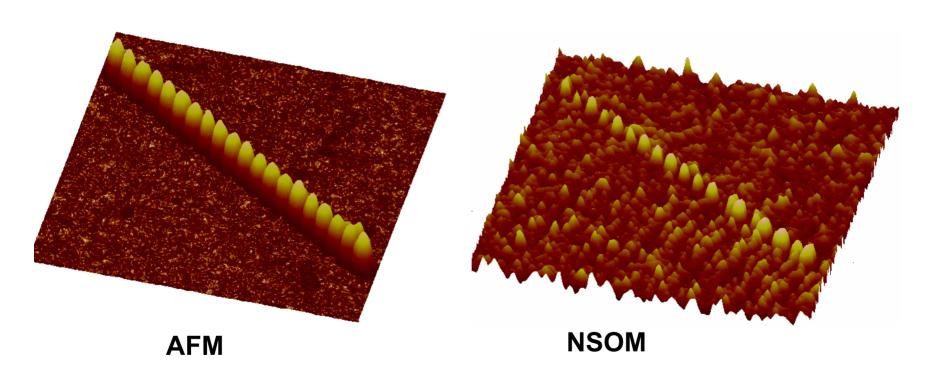
Lots of interesting experiments are being carried out throughout the world -- theory and computation are needed!



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Chain of Ag Nanoparticles

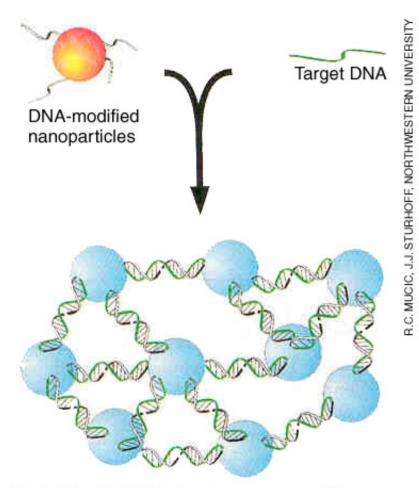
Wiederrecht and co-workers (ANL): AFM (topography) and NSOM (near field strengths) on Ag particle arrays



size ~ 100 nm width, periodicity ~ 200 nm





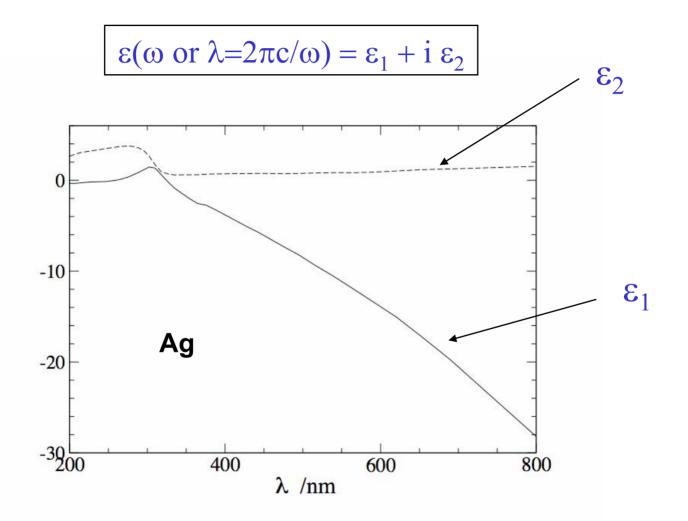


Hybrid web. DNA strands on nanoparticles seek out target DNA, forming a network.

Most potential applications aim to tap the properties of surface plasmons (SPs) -- resonant interactions of light with (induced) charge density in the MNP surface.

E.g., in Mirkin's biological sensors, the SP resonance scattering blue-shifts when Au MNPs bind to DNA in manner shown.

Origins of Surface Plasmons: $Re[\varepsilon] < 0$ in Metals



Standard heuristic argument

E.o.m. for electron (in a metal) displacement in an electric field:

$$mXY + (m/\tau) \dot{X} = -q E_x e^{-i\omega t}$$

Assume
$$x(t) = x_0 e^{-i\omega t}$$
 to get

$$x_o = (q E_x/m)/(\omega^2 + i\omega/\tau)$$

Polarization for number density *n*:

$$P_{x} = n (-qx_{o}) = \chi \varepsilon_{o} E_{x}$$

 $= (\varepsilon - 1) \varepsilon_0 E_x$

Solve for relative dielectric constant, ε :

$$\varepsilon(\omega) = 1 - \omega_p^2/(\omega^2 + i\omega/\tau)$$

$$\omega_p = \sqrt{nq^2/m\varepsilon_o}$$
 Drude model





Surface Plasmons

For $Re[\varepsilon]$ < 0, strong oscillations in electric charge are predicted which allow for electromagnetic surface waves, which can have longitudinal components but decay rapidly away from a surface ("evanescent")

Local Surface Plasmons (LSP's): Confined to nanoparticle surfaces or nanohole inner surfaces; akin to bound (or resonance) states in QM.

Surface Plasmon Polaritons (SPP's): Confined to metal surfaces but traveling surface plasmons. Typically generated on flat thin metal films.



Methods

"Real theory", if possible, is best but for exploring somewhat complicated nanophotonics systems robust, reasonably accurate and easy to use computational methods are also helpful.



Finite-Difference Time-Domain (FDTD) Method

Basic idea and numerical method is somewhat old:

K. S. Yee, IEEE Trans. Antennas and Propagation, **14**, 302 (1966).

Modern implementations are discussed in:

A Taflove and S. C. Hagness, *Computational Electrodynamics: The Finite-Difference Time-Domain Method*, 2nd Ed., (Artech House, Boston, 2000).





- Solve t-dependent form of Maxwell's equations: discretizing in both space and time to generate **E**(**x**, t) and **H**(**x**, t) via straightforward time-stepping.
- Staggered grids used for more accurate finite differences and preservation of certain symmetries.
- Current fields are introduced to describe the metallic dielectric constant behavior which can have $\text{Re}[\epsilon(\omega)] < 0$.



Maxwell's Equations

Outside each nanoparticle:

$$\partial \mathbf{E}(\mathbf{x},t)/\partial t = \nabla \times \mathbf{H}(\mathbf{x},t)/\varepsilon(\mathbf{x})$$

$$\partial \mathbf{H}(\mathbf{x},t)/\partial t = -\nabla \times \mathbf{E}(\mathbf{x},t)/\mu_{o}$$

Inside each nanoparticle region:

$$\partial \mathbf{E}(\mathbf{x},t)/\partial t = [\nabla \mathbf{x} \mathbf{H}(\mathbf{x},t) - \mathbf{J}(\mathbf{x},t)]/\varepsilon_{\infty}$$

$$\partial \mathbf{H}(\mathbf{x},t)/\partial t = -\nabla \times \mathbf{E}(\mathbf{x},t)/\mu_{o}$$

$$\partial \mathbf{J}(\mathbf{x},t)/\partial t = \epsilon_o \omega_p^2 \mathbf{E}(\mathbf{x},t)/\mu_o - \nu \mathbf{J}(\mathbf{x},t)$$

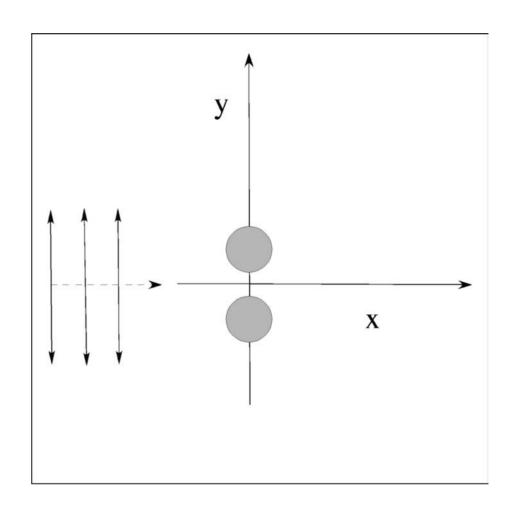
(Parameters for current terms chosen to fit metal dielectric constant data in some frequency range)





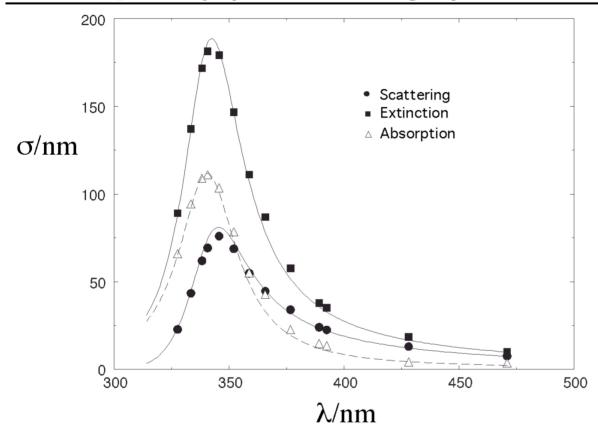
Metal Nanowires

Nanoscale cross sections with infinite extension out of (x,y) plane:





FDTD approach is reasonably quantitative even for frequency (or wavelength) resolved scattering



Single Ag cylinder, radius 25 nm:

SP resonance at 350 nm

Symbols: FDTD result based on Fourier transforms of numerically generated, time-dependent fields Curves: Analytical results ("Mie theory").

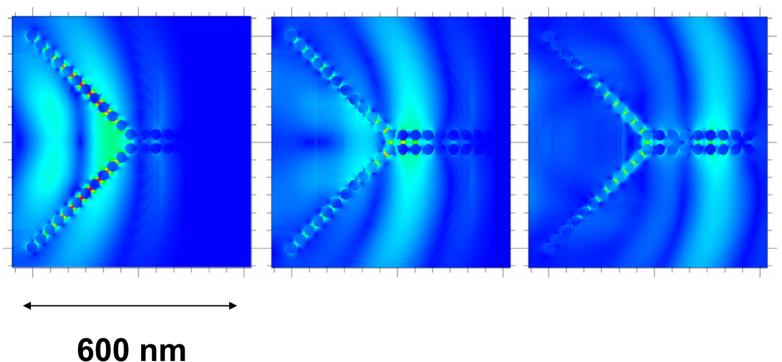




Funnel Configuration of Ag Nanowires

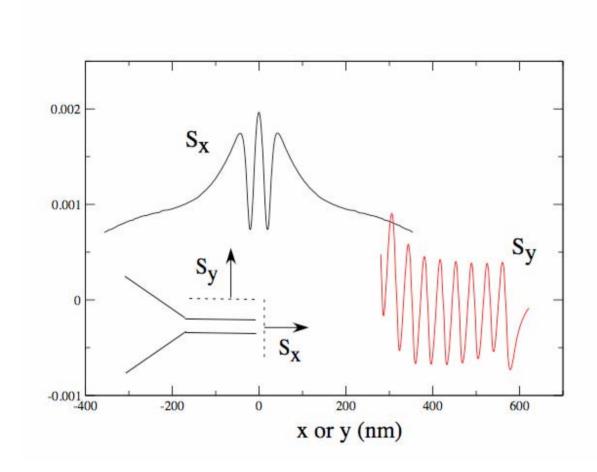
[Gray and Kupka, Phys. Rev. B 68, 045415 (2003)]

E field at 0.5 fs time intervals:



Can achieve 100 nm nanoscale localization of light.

Time-averaged Flux



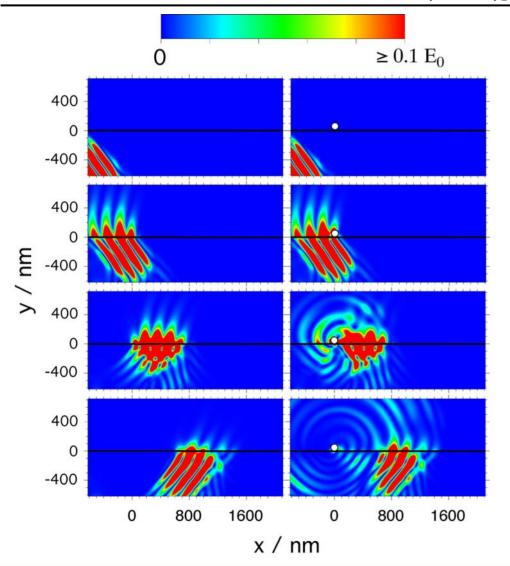


Evanescent Excitation and Scattering of Ag Nanowires

- Wurtz, Im, and Wiederrecht: find metal nanoparticles, scatter light at small angles above dielectric/air interface in "Total Internal Reflection" experiments.
- Calculations have begun to explore this process.



TIR Excitation with FDTD [Wurtz et al., J. Phys. Chem. B. **51**, 14191 (2003)]



Left panels, top to bottom, show TIR in absence of nanowire.

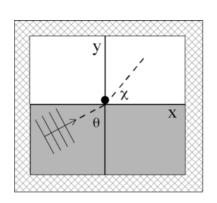
Right panels, top to bottom, show what happens with a 50 nm radius Ag cylinder.

Glass (n=1.5) for y < 0, air (n=1) for $y \ge 0$.



Low angle scattering

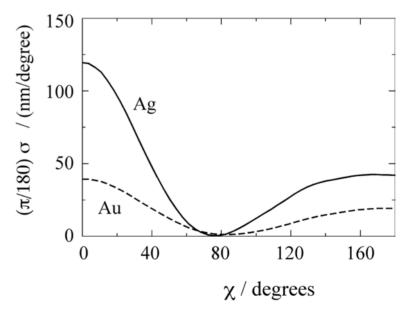
Incident light closer to Ag SP resonance than the Au one.



1 particle/surface

 χ / degrees

1 particle, free space



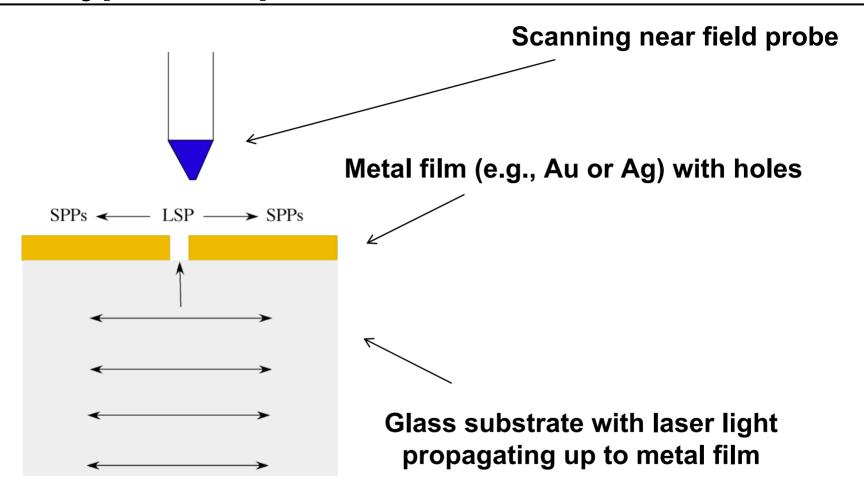


Nanoholes in Thin Metal Films (S.-H. Chang)

- Metal films with nanoscale holes are also of interest -- local SPs similar to those in metal nanoparticles can be excited on hole edges. Furthermore, coupling to SPP's can also occur.
- L. Yin, U. Welp, V. Vlasko-Vlasov and coworkers are exploring such nanoholes at Argonne.



The Type of Experiment:

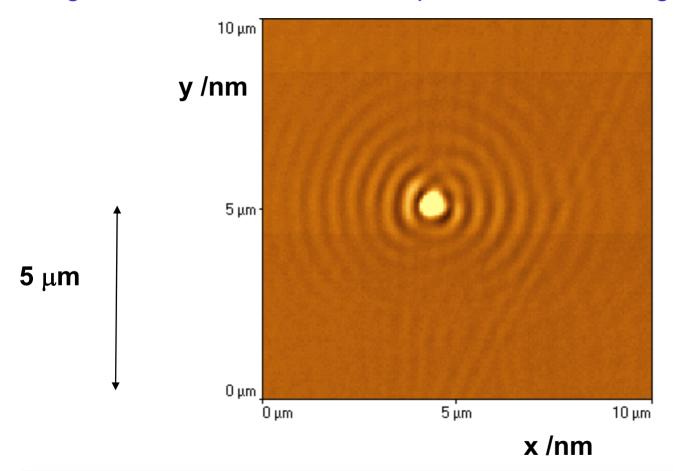






Experimental Results: \(\lambda \) Fringes

170 nm diameter hole in a 100 nm thick Au film on glass. 530 nm light. A view from the air-Au top: \approx 470-480 nm fringes



x-polarized light



$$\lambda_{SPP} = \sqrt{\frac{\mathcal{E}_{M} + 1}{\mathcal{E}_{M}}} \lambda_{inc} = \sqrt{\frac{-4.67 + 1}{-4.67}} 532 nm \approx 472 nm$$

But SPPs alone cannot be seen easily:

$$E_x \approx C \exp(ik_{SPP}x)$$

 $\Rightarrow I = E_x E_x^* \approx C^2$

Suppose instead:

$$E_x \approx C \exp(ik_{SPP}x) + A \exp(iS)$$

Then:

$$I = E_x E_x^* \approx C^2 + A^2 + 2CA \cos(k_{SPP}x + S)$$

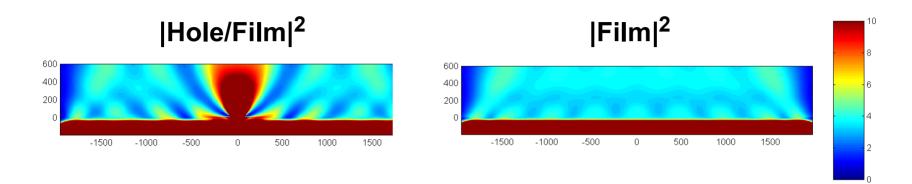
$$k_{SPP} = 2\pi/\lambda_{SPP} \Rightarrow \lambda_{SPP}$$
 fringes

But what is A exp(iS)?

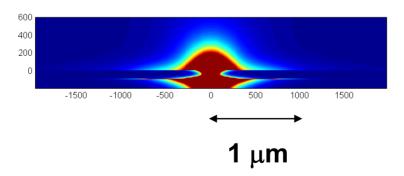




$A \exp(iS) = T = a Directly Transmitted Wave$



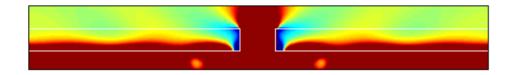
| Hole/Film - Pure Film|²



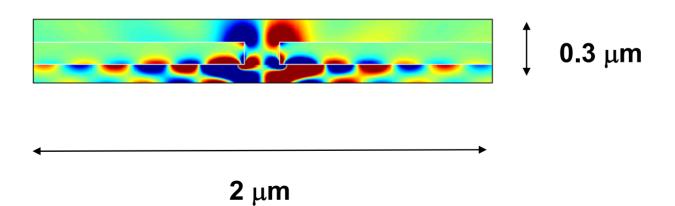


Poynting Vector

Flux upwards



Flux sideways



Nanoholes act as "Point Sources" of SPP's

Near the top plane of the metal surface, we find:

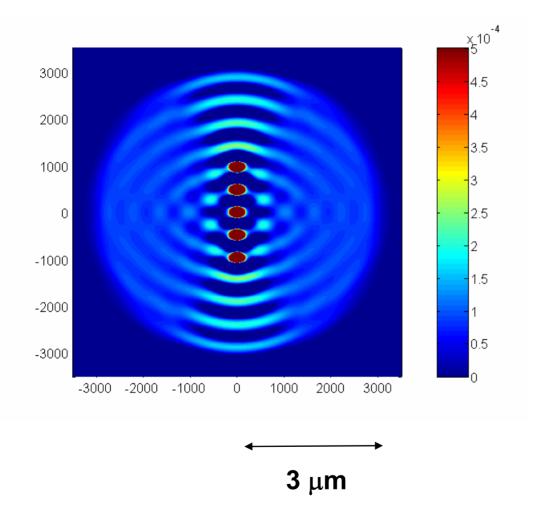
$$E_x(r,\varphi,z) \propto C \frac{\exp(ik_{SPP} r - \gamma z)}{\sqrt{r}} \cos(\varphi) + T_x$$

where x-polarized $[\cos(\varphi)]$ light is assumed.

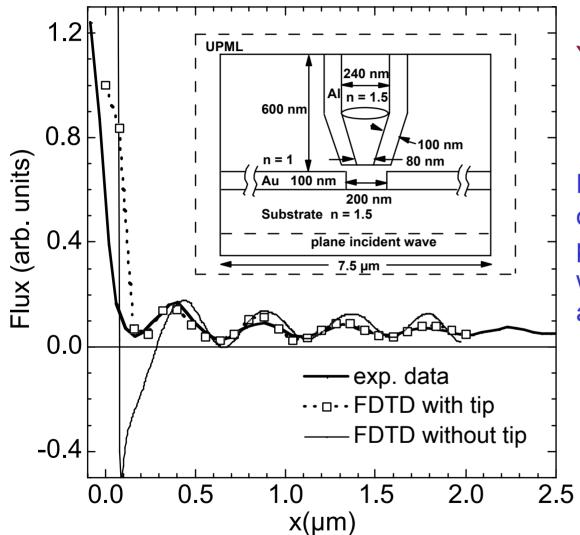
• Despite fact λ_{inc} , $\lambda_{SPP} > 3$ d_{hole}, a limited kind of "Huygen's principle" appears to apply to planar SPP's ...



Nanohole Arrays for Stronger SPP Beams



Can also explicitly model the NSOM probe

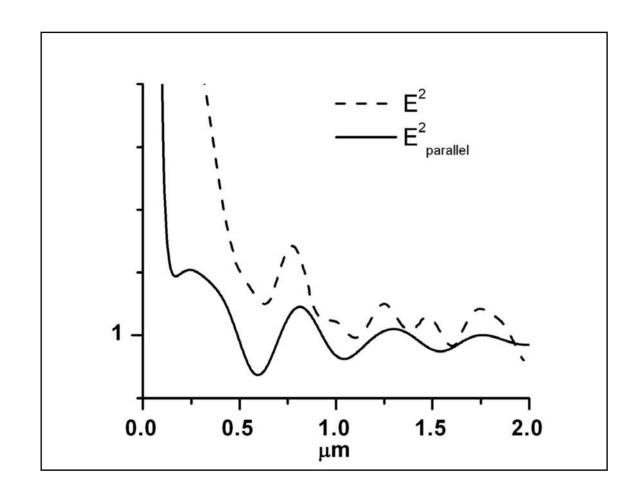


Yin et al., *Appl. Phys. Lett.* **85**, 467 (2004).

Probe develops dipolar oscillation with incident (x) polarization, interacts with/picks up information about $|E_x|^2$ (not $|E_{tot}|^2$).



Parallel (x) component shows more clear fringes







Concluding Remarks

- Demonstrated FDTD approach can yield optical cross sections for metallic nanowire problems.
- Arrays of nanowires investigated: funnel configurations show 100 nm scale localization of light.
- Total Internal Reflection (TIR)
 excitation/scattering by nanoparticle systems investigated.
- Nanoholes investigated -- demonstrated how SPP's reveal themselves through interference.





Future Directions

- Multiscale Modeling of Nanophotonics
- Coherent Control of Nanophotonics
- "Killer Apps" ??



Multiscale Approach to Modeling Nanophotonics Problems:

